Impedance sensing and control

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Overview

- Cavity impedance condition
- Impedance readout
- Experimental demo of IM control
- Applications of IM control
- Impedance control for (Ad)Virgo



Impedance matching of resonant optical systems



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Impedance matching of resonant optical systems

- Impedance matching for optically resonant systems is a kin to that in electrical systems:
 - In electrical systems impedance matching provides optimum voltage, power, or current transfer.
 - In optical systems impedance matching provides the optimum electric field transfer.



Impedance matching of resonant optical systems

- Impedance matching for optically resonant systems is a kin to that in electrical systems:
 - In electrical systems impedance matching provides optimum voltage, power, or current transfer.
 - In optical systems impedance matching provides the optimum electric field transfer.
- Interrogation and control of the impedance matching condition offers an alternative active feedback control technique that has applications within GW interferometry, absorption spectroscopy, and quantum optics experiments.



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The important parameters that active impedance matching optimises are:

• optimum electric field transfer through the optical system.



The important parameters that active impedance matching optimises are:

- optimum electric field transfer through the optical system.
- ensures that the reflected electric field is zero and the circulating field is maximised (assuming a back mirror who's reflectivity is dominated by loss).
- by optimising the circulating power, the technique also optimises the signal sensitivity of the system.











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• By monitoring the beat signal between the carrier and a set of amplitude modulated sidebands which are well outside the coupled cavity linewidth, we can obtain a signal which is proportional to the reflected electric field amplitude:

$$E_{refl} \approx E_{inc} e^{i\omega_c t} \left[\mathcal{R}(\omega_c) + \frac{\beta}{2} \mathcal{R}(\omega_c + \omega_m) e^{i\omega_m t} + \frac{\beta}{2} \mathcal{R}(\omega_c - \omega_m) e^{-i\omega_m t} \right]$$

• Demodulation allows to extract a signal which is linearly dependent on the reflected amplitude response of the coupled cavity.

$$\mathcal{E}_{\mathcal{Q}} = \sqrt{P_c P_s} \times Re[\mathcal{R}(\omega_c)\mathcal{R}^*(\omega_c + \omega_m) + \mathcal{R}^*(\omega_c)\mathcal{R}(\omega_c - \omega_m)].$$



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• The photo detector signal is:

$$P_{refl} = P_c |\mathcal{R}(\omega_c)|^2 + P_s |\mathcal{R}(\omega_c + \omega_m)|^2 + P_s |\mathcal{R}(\omega_c - \omega_m)|^2 + 2\sqrt{P_c P_s} \times Re[\mathcal{R}(\omega_c)\mathcal{R}^*(\omega_c + \omega_m) + \mathcal{R}^*(\omega_c)\mathcal{R}(\omega_c - \omega_m)]\cos(\omega_m t) + 2\sqrt{P_c P_s} \times Im[\mathcal{R}(\omega_c)\mathcal{R}^*(\omega_c + \omega_m) + \mathcal{R}^*(\omega_c)\mathcal{R}(\omega_c - \omega_m)]\sin(\omega_m t) - 2P_s\cos(2\omega_m t).$$

$$(6.2)$$

• Subsequent demodulation allows to extract a signal which is linearly dependent on the reflected amplitude response of the coupled cavity.

$$\mathcal{E}_{\mathcal{Q}} = \sqrt{P_c P_s} \times Re[\mathcal{R}(\omega_c)\mathcal{R}^*(\omega_c + \omega_m) + \mathcal{R}^*(\omega_c)\mathcal{R}(\omega_c - \omega_m)].$$



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Changing the impedance



- We can use a Michelson as a variable reflectivity mirror.
- Or we can use a Fabry-Perot cavity as a variable reflectivity mirror.



Experimental PD and error signals





Experimental PD and error signals



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Impedance error signal



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Frequency response data



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PDH sensitivity VS impedance





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Absorption spectroscopy







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Metrology for evanescent coupling in micro-ring cavity





